Aspects of transition metal pentacoordination: investigation of the structure and bonding of Me₃NbCl₂, Me₃TaCl₂ and Me₂NbCl₃ by photoelectron spectroscopy and density functional theory calculations

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He I and He II photoelectron (PE) spectra of Me₃NbCl₂, Me₂NbCl₃ and Me₃TaCl₂ are presented and assigned. Density functional calculations give trigonal bipyramidal structures for the Nb complexes, with the axial positions occupied by Cl ligands. Calculated ionisation energies (IE) for Me₃NbCl₂, Me₂NbCl₃ and Me₃TaCl₂ are in good agreement with the experimental values and the molecular orbital compositions can be reconciled with the intensity changes in the PE spectra. The axial interactions are best described by a 3-centre 4-electron bond, which accounts for the longer Nb–Cl distances predicted for these bond lengths. That the orbitals involved in Ta–ligand bonding with a Ta 6s contribution have higher IE's relative to their Nb counterparts may be due to relativistic stabilisation of the Ta 6s orbital. Comparisons with the main group analogue Me₃SbCl₂ support Elbel's assignment of the PE spectrum of the latter compound, where there is a reversal of the IE ordering between the transition metal and the main group compounds. For Sb the halide orbitals constitute the highest lying set, whereas with the transition metal compounds the metal–carbon bonding orbitals provide the highest occupied molecular orbitals. No evidence is found in the calculations for a significant contribution to the bonding from Nb 5p orbitals. The large angles subtended by the Cl atoms at the metal are attributed to their negative charge and interligand repulsion.

Introduction

Five-coordinate species are remarkable in structural chemistry, as they have no low energy geometry with stereochemically equivalent sites. As such, pentacoordinate main group compounds played a central role in the development of VSEPR theory in the 1960's.^{1,2} The trigonal bipyramid (tbp) was found to be favoured over the square pyramid (sp) in most cases, and the concept of apicophilicity was introduced to account for ligand site preferences in heteroleptic tbp systems.^{3,4} Fivecoordinate transition-metal centres figure prominently in many reaction pathways, as well as featuring in ground state geometries. In 1975 a seminal paper by Rossi and Hoffmann⁵ employed extended Hückel calculations to address the bonding in five-coordinate transition-metal complexes. The absence of non-bonding d electrons in the mixed ligand species Me₃NbCl₂, Me₃TaCl₂ and Me₂NbCl₃ has prompted us to study these systems by a variety of techniques. Here we report our investigation of the bonding by a combination of photoelectron spectroscopy and density functional theory (DFT) calculations. The He I spectrum of Me₃TaCl₂ has been reported previously;⁶ here we report the He II spectrum which suggests a minor change in the earlier assignment of the bands.

Experimental

Synthesis and characterisation

Standard high-vacuum and inert atmosphere techniques were used throughout. Me_2NbCl_3 , Me_3NbCl_2 and Me_3TaCl_2 were prepared by the reaction of NbCl₅ or TaCl₅ with Me_2Zn in

pentane according to the procedures of Juvinall and Fowles *et al.*^{7,8} The volatile materials were separated by fractional sublimation *in vacuo*. Sample purity was checked by reference to the ¹H NMR spectrum of a CD₂Cl₂ solution,^{8,9} and to the IR spectrum of the solid condensate at 77 K.¹⁰ Samples were stored at 77 K until needed to prevent thermal decomposition.

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Photoelectron spectroscopy

He I and He II PE spectra of Me_3NbCl_2 , Me_3TaCl_2 and Me_2-NbCl_3 were measured using a PES Laboratories 0078 spectrometer interfaced with an Atari microprocessor, which enables accumulation of a spectrum by repeated scans. Spectra were calibrated using He, Xe and N₂.

Computational methods

All calculations were carried out using the Amsterdam density functional (ADF) program system, versions 2.0.1 and 2.3.¹¹ For geometry optimisations and calculation of ionisation energies, the electronic configurations of the molecular systems were described by an uncontracted triple- ζ basis set of Slater-type orbitals (STO). Hydrogen, carbon and chlorine were given an extra polarisation function, 2p on H and 3d on C and Cl. The cores of the atoms were frozen, C up to 1s, Cl to 2p, Nb up to 3d, Sb up to 4p and Ta up to 5p. First order relativistic corrections were made to the cores of all atoms using the Pauli formalism.

Energies were calculated using Vosko, Wilk and Nusair's local exchange correlation potential¹² with non-local-exchange corrections by Becke¹³ and non-local correlation corrections by Perdew.^{14,15} The non-local correction terms were utilised in calculating gradients during geometry optimisation, so as to find the non-local minimum.

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Table 1 Calculated bond lengths and angles for Me_3ECl_2 (E = Nb, Ta and Sb) with C_{3v} symmetry and Me_2NbCl_3 with C_{2v} symmetry

	M- NI-CI		Me ₃ TaCl ₂		Me ₃ SbCl ₂	
Parameter	Calc.	Calc.	Calc.	Exp. ^a	Calc.	Exp. ^{<i>b</i>}
r(M–Cl) _{ea} /Å	2.30					
$r(M-Cl)_{av}/Å$	2.34	2.35	2.41	2.317(3)	2.53	2.460(6)
r(M-C)/Å	2.17	2.17	2.20	2.158(5)	2.15	2.107(6)
$r(C-H)_{av}/Å$	1.10	1.10	1.10	1.098(7)	1.10	1.103(13)
Clea-M-Clax/°	98					
Clax-M-Clax/°	164	180 ^c	180 ^c	180 ^c	180 ^c	180 °
Clea-M-C/°	121					
Cl _{ax} -M-C/°	86	88, 92	88, 92	90 ^c	89, 91	90 ^c
C–M–C/°	119	120 ^c	120 ^c	120 ^c	120 ^c	120 °

 a^{a} r_{a} Structure determined by electron diffraction studies on the gaseous molecule, see ref. 16. b^{b} r_{a} Structure determined by electron diffraction studies on the gaseous molecule, see ref. 17. c^{c} Fixed by symmetry constraint.



Fig. 1 Initial structures assumed in geometry optimisations of Me₃-NbCl₂ and Me₂NbCl₃.

Ionisation energies were estimated by calculating energies for the molecular ions in their ground and excited states. In all cases the geometries were fixed at that found for the optimum structure of the molecule to enable comparison with the experimental vertical IE. A complication arose in the case of Me₃ECl₂ (E = Nb, Sb and Ta), as the ADF program used does not handle C_{3h} symmetry explicitly. This made the calculation of the excited states of the molecular ion unfeasible with this particular geometry constraint. The states were thus calculated using C_{3v} symmetry, which involved rotation of the methyl groups through 30° from the C_{3h} geometry so that one C–H bond in each CH₃ group eclipsed the same M–Cl bond. Optimum bond lengths were found to be very close to those of the C_{3h} structure, as were the corresponding one-electron energies.

For the calculation of ion states the basis set used was the molecular orbital (MO) from the molecular calculation. This assisted convergence on a hole in the specified orbital, and also provided a ready check on the extent to which the ion state hole could be described in terms of a single MO of the parent molecule.

Results

Structural predictions

For the Nb compounds under study, geometry optimisations were carried out with initial structures based on a tbp framework with various site occupancies; these are shown in Fig. 1.

Use of symmetry constraints enabled an estimation of the relative energies of the various conformers envisaged for Me₂NbCl₃. The most stable was found to be a C_{2v} structure for **2a**, confirming the relative preferences of Cl for axial and Me for equatorial sites. The D_{3h} conformer lay 1.43 eV in energy above the C_{2v} structure. The C_s conformer optimised to a geometry where the angle subtended by the original axial groups at the Nb was 155°; this was 0.33 eV higher in energy than the C_{2v} model. A frequency calculation confirmed that the C_{2v} structure was a local minimum. The lowest energy vibrational mode, with a calculated wavenumber of 80 cm^{-1} , involved bending of the two axial Cl and the equatorial Me ligands towards an sp structure, the process involved in Berry pseudorotation.

In the case of Me₃NbCl₂ the C_3 structures had the lowest energy, the difference between the C_{3h} and C_{3v} forms being negligible. Structure **1b** moved to an sp form with one Cl ligand axial, and having an energy 0.33 eV above those of the C_3 conformers. Model **1c** moved to a distorted tbp geometry with an angle of 148° between the axial Me and 126° between the equatorial Cl ligands; this lay 0.48 eV above the C_3 conformers. Frequency calculations on the optimised C_{3h} and C_{3v} structures gave three very low imaginary frequencies for both, corresponding to rotation of the methyl groups. However, optimisation with no symmetry also failed to locate an energy minimum for the relative orientation of the methyl groups.

The predicted bond lengths and angles for the optimised structures of Me_3NbCl_2 and Me_2NbCl_3 are given in Table 1. Although the angle of 164° subtended by the two axial Cl's at the Nb in Me_2NbCl_3 is appreciably less than the ideal of 180°, the skeletal shape is clearly closer to a tbp than an sp. The two axial Cl's bend away from the equatorial Cl. The Nb–Cl axial bonds are calculated to be longer than the equatorial Nb–Cl, indicating that Cl is more strongly bound in an equatorial site. Very small distortions of the Me groups are predicted but the variations in bond length and angle are likely to be beyond the limits of experimental detection.

For comparative purposes, geometry optimisations assuming C_{3h} and C_{3v} symmetry were also carried out on Me₃TaCl₂ and Me₃SbCl₂, with the results also given in Table 1. In these two cases, comparisons with the structures determined previously by electron diffraction studies of the vapours^{16,17} find satisfactory agreement between theory and experiment.

Photoelectron spectroscopy

He I and He II spectra of Me₃NbCl₂, Me₃TaCl₂ and Me₂NbCl₃ are shown in Fig. 2-4. Vertical ionisation energies are tabulated in Table 2. The He I spectrum recorded for Me₃TaCl₂ is in good agreement with that previously published.6 For all three compounds, there are substantial differences between the band intensities generated with the different photon energies. A previous study of the PE spectrum of MeTiCl₃, with photon energies varying between 22 and 50 eV, showed similar variations in photoionisation cross-sections.¹⁸ Such cross-section variations are important in their capacity to give a clear indication of the atomic orbitals (AO) contributing to the molecular orbital with which a particular PE band is associated. For example, Nb 4d ionisations will gain in relative intensity on increasing the photon energy. Cl 3p ionisations will fall dramatically in relative intensity between He I and He II radiation, as Cl 3p has a Cooper minimum in its cross-section around 42 eV. C 2p cross-sections show an intermediate behaviour. Thus,



Fig. 2 He I and He II PE spectra of Me₃NbCl₂.



Fig. 3 He I and He II PE spectra of Me₃TaCl₂.

in the He I spectrum the Cl 3p ionisations will dominate the spectrum, whereas at 40 eV Nb 4d will be the dominant ionisation.

The IE of a band is also an indication of its orbital character. By analogy with MeTiCl₃ and other metal halides the lowest IE PE bands will be due to metal–carbon bonding electrons, followed by bands associated with Cl lone pair $p\pi$ ionisations, followed by metal chlorine bonding ionisations, with C–H bonding ionisations having the highest IE in the valence region.¹⁸

For Me_3NbCl_2 (Fig. 2), bands A, B and D come from orbitals with metal character and C from orbitals with Cl character. Bands A and B can be assigned to ionisation of Nb–C bonding electrons, band C to Cl p ionisations, and band D to Nb–Cl



Fig. 4 He I and He II PE spectra of Me₂NbCl₃.

bonding ionisations. The broad band E, which to a certain extent underlies D, is characteristic of the ionisation of C–H bonding electrons.

The PE spectrum of Me_3TaCl_2 (Fig. 3) is similar to that of its Nb analogue, if it is assumed that the equivalent of band B forms the low IE part of the second complex band. The intensity variations between the He I and He II spectra show a change in band shape and an increase in the intensity of the lower IE part of the band consistent with its association with electrons having metal character. In this case, band E is associated with an orbital of metal character superimposed on a broad C–H band D.

The PE spectrum of Me₂NbCl₃ may be assigned in a similar fashion. Bands A and B are associated with Nb–C bonding orbitals, C and D with Cl 3p ionisations, and E and F with Nb–Cl bonding electrons. The broad band G is due to C–H bonding ionisations.

For all three compounds, a more detailed assignment is made possible by examination of the DFT bonding model and calculation of the relevant IE.

MO analysis

In the MO analysis we are principally concerned with the framework orbitals and the Cl $p\pi$ electrons. For the Me₃MCl₂ series, the framework has effective D_{3h} symmetry which distinguishes some orbital combinations otherwise possessing identical irreducible representations in the lower C_3 symmetry. Accordingly, we use both labels to describe the orbitals.

Pictorial representations of selected one-electron wave functions of Me₃NbCl₂ are given in Fig. 5. One-electron energies and a Mulliken population analysis of the upper occupied orbitals are detailed in Table 2. Orbitals 7e (e') and 9a₁ (a₁') show metal–carbon bonding characteristics, the orbital 9a₁ (a₁') being formed with a Nb (s–d₂2) hybrid. They are followed, on an increasing energy scale, by a set of orbitals of high chlorine 3p character, 6e (e''), 5e (e') and 8a₁ (a₂''). The 5e (e'') orbital shows some Nb character, which is π bonding with respect to the axial Cl. The 8a₁ (a₂'') orbital has Cl p σ character; in principle this symmetry adapted combination could form a bonding interaction with a Nb 5p orbital, but the Mulliken analysis and the energy suggest its bonding role to be minimal. The 7a₁ (a₁') orbital represents bonding between the axial chlorine p σ

Table 2	Calculated orbital energies,	calculated IE, experimental	vertical IE and principal orbital	character for Me ₃ MCl ₂ molecules ^{<i>a</i>}
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					Population analysis			
Orbital	Orbital	IE/eV			Nb			
	Orbital energy/eV	Calc.	Calc. Exp.	Orbital character	d	s	p	p p
7e, e'	-6.99	9.64	9.85 A ₁	$xy, x^2 - y^2 + C$	27		9	41
Qa a '	-7.83	10.45	10.14 A ₂	$s = r^2 + C \sigma$	15	12		54
a_1, a_1	-7.02	10.45	10.6 D	S = 2 + C O	0	12	74	54
50, C	-8.01	11.12	11.05 C	$Cl \pi$	9		7 4 86	
Je, e 80 0 "	-8.01	10.06	11.05 C	$Cl\pi$			02	
a_1, a_2	-0.77	12.40	12.05 C				92	
$7a_{1}^{1}, a_{1}^{2}$	-9.85	12.59	13.4 D	$s + z^2 + Cl \sigma$	14	3	66	
(b) M = Ta	1							
					Popul	lysis		
		IE/eV			Та			
	Orbital						Cl	С
Orbital	energy/eV	Calc.	Exp.	Orbital character	d	s	р	р
6e, e'	-6.93	9.59	9.91 A ₁ 10.26 A ₂	$xy, x^2 - y^2 + C \sigma$	27		6	49
5e. e″	-7.89	10.61	11.42 B	Cl π	5		79	9
4e. e'	-7.96	11.10	11.85 C	$Cl \pi$	5		79	5
7a. a."	-8.01	10.89	11.42 C	Clo			96	
6a ₁ , a ₁ '	-8.32	10.97	11.85 B	$s - z^2 + C \sigma$	20	12		50
1a ₂ , a ₂ '	-9.85	12.49	12.8–15.5 D	C–H				
$4a_1, a_1'$	-10.17	13.03	13.85 E	$s + z^2 + Cl \sigma$	10	9	69	3
(c) M = Sb)							
					Population analysis			
		IE/eV		Principal	Sb		~	~
Onlyit-1	Orbital	Cal-	Euro ^b	anhital abarrates			CI	C
Orbitai	energy/eV	Calc.	Exp.	orbital character	s	р	р	p
8e, e'	-6.72	9.48	9.8	Cl π			91	5
7e, e″	-6.84	10.14	10.22	Cl π			96	
$8a_1, a_1'$	-7.09	9.81	10.65	Cl σ			83	11
$7a_1, a_2''$	-8.74	11.41	12.05	$Cl \sigma + Sb p$		15	63	6
6e, e'	-8.91	11.60	12.05	$C \sigma + Sb p$		22		49
1a ₂ , a ₂ '	-9.85	13.33		C–H Î				

" The orbital numbering differs between Nb, Ta and Sb because the Nb calculations include the Nb 4p core orbitals, while the Ta calculations treat the Ta 5p core orbitals as frozen, and the Sb calculations include the 4d orbitals in the valence set, while the 3p orbitals are frozen. ^b See ref. 22.

orbitals and the Nb d_{z^2} orbital, chlorine character predominating. The Nb–Cl σ bonding is thus consistent with the simple model of a 3-centre 4-electron bond, but with a high degree of polarity. The 1a₂ orbital is constrained to be purely C–H in character and is a convenient marker for the onset of the C–H bonding manifold.

Evaluation of the ionisation energies associated with removing electrons from these orbitals (Table 2) enables the assignment of the PE spectrum of Me₃NbCl₂ to be completed. The correlation between the calculated and experimental IE, if not obvious at first, is made clear by consideration of the intensity changes and the Mulliken populations. Band A is assigned to ionisation from the 7e orbital. The fact that it shows two maxima, A₁ and A₂, suggests that the ²E ground state of the molecular ion undergoes Jahn–Teller distortion. Band B is associated with ionisation from the 9a₁ orbital. The very low intensity of this band in the He I spectrum, such that it is barely visible, is consistent with the intensity patterns found from the ionisation of other transition metal halides with orbitals having significant metal s character.¹⁸ Band C is associated with the orbitals of high Cl character, *i.e.* 6e, 5e and $8a_1$. Band D is attributed to the $7a_1$ ionisation. The bonding character of the orbital ensures the IE is higher than for ionisation of the Cl non-bonding electrons, while its metal d character has the consequence that band D grows relative to band C as the ionising radiation is changed from He I to He II. Calculation of the IE from an upper C–H bonding orbital places the onset of the C–H ionisation manifold very close to band E, in support of the preliminary interpretation of this band.

 Me_3TaCl_2 has an orbital structure very similar to that of its Nb analogue. Accordingly, the PE spectrum is given a parallel assignment. The principal difference from Me_3NbCl_2 is revealed in the higher IE for the a_1' framework orbitals. The Ta-C a_1' bonding orbital is stabilised to an extent that places its ionisation (band B) in the second complex band overlapping with the Cl p bands. In addition, the Ta-Cl a_1' IE is greater so that the associated band E lies around the maximum of the C-H band manifold and does not precede it, as in the case of Me_3NbCl_2 . As both these orbitals have Ta 6s content, their greater binding energies may well be a consequence of the

Table 3 Calculated orbital energies, calculated IE, experimental vertical IE and principal orbital character for the Me₃NbCl₃ molecule

					Population analysis			
		IE/eV			Nb		C1	
Orbita	l orbital energy/eV	Calc.	Exp.	Orbital character	d	8	Cl p	C p
3b ₂	-7.05	9.64	9.89 A	$xz + C \sigma$	21		32	36
4a ₁	-7.73	10.35	10.64 B	$x^2 + C \sigma$	19	4	26	29
3b ₁	-7.81	10.53	11.0 C	Cl π			99	
1a,	-8.09	10.75	11.0 C	$xy + Cl \pi$	9		70	8
3a ₁	-8.27	11.00	11.64 D	Ċl π	6		68	9
2b ₂	-8.32	11.00	11.64 D	Cl π			83	7
2b ₁	-8.54	11.29	11.64 D	Clσ			84	4
1b,	-8.77	11.09	11.64 D	$xz + Cl \pi$	16		73	6
1b ₁	-9.23	12.00	12.73 E	$yz + Cl \pi$	16		74	3
$2a_1$	-9.77	12.52	13.57 F	$z^2 + Cl \sigma$	8		71	8
1a ₁	-10.18	12.56	13.57 F	$y^2 + Cl \sigma$	18		73	





Fig. 5 Selected representations of the MO's of Me₃NbCl₂.

relativistic stabilisation of the Ta 6s orbital.¹⁹⁻²¹ Elbel et al. did not identify any ionisation from the Ta-C a₁ ionisation orbital,⁶ but otherwise their assignments correspond with ours.

For Me₂NbCl₃, the one-electron wave functions are represented pictorially in Fig. 6. One-electron energies and a Mulliken population analysis of the upper occupied orbitals are given in Table 3. Highest in energy in the occupied orbital energy manifold come the $3b_2$ and $4a_1$ orbitals, which have metal-carbon bonding character (Fig. 6). Of the next five orbitals, which are principally localised on Cl, the 1a2 orbital shows some Nb–Cl π bonding character (Fig. 6). The 2b₁ orbital, which is Cl $p\sigma$ in character and located on the axial ligands, has a negligible Nb 5p contribution, being very similar in its distribution to the a₂" orbital of Me₃NbCl₂. The 1b₁ and $1b_2$ orbitals are also Nb-Cl π bonding in character. Finally, the $2a_1$ and $1a_1$ orbitals provide the Nb-Cl σ bonding.

The greater number of ionisations of interest for Me₂NbCl₃ as compared with Me₃NbCl₂ and the lower symmetry of Me₂NbCl₃, hamper the calculation of the ion state energies, as there are now several with the same symmetry representations.



Fig. 6 Selected representations of the MO's of Me₂NbCl₃.

However, the calculations do give valuable guidance in assigning the PE spectrum (Table 3). Bands A and B are ascribed to ionisation from the 3b₂ and 4a₁ orbitals, respectively. Again the second band B is hard to see in the He I spectrum. Even though the orbital is of a_1 symmetry, the population analysis does not show a 5s contribution as high as that found for the analogous a₁' orbital of Me₃NbCl₂. On the evidence of the IE pattern calculated and the relative intensities of the bands, it seems reasonable to assign band C to orbitals 3b1 and 1a2 and band D to 3a₁, 2b₂, 2b₁ and 1b₂, as these all have high Cl content and the associated ion states are predicted to lie within 0.3 eV of one another. Working again from the ion state energy pattern leads us to assign band E to the ionisation from the 1b₁ orbital and band F to the $2a_1$ and $1a_1$ ionisations.

Discussion

Comparison with transition metal analogues

The appearance and assignment of the spectrum of Me₃NbCl₂ accord well with that of its heavier congener Me₃TaCl₂, which was studied by Elbel et al.⁶ Correlation of the orbital energies is shown in Fig. 7, where the most notable feature is the stabilisation of the M-C and M-Cl a1' orbitals on descending the group. This correlates with the increased strength of the bonds

Table 4 Mulliken charges calculated for the molecules Me_3MCl_2 (M = Nb, Ta or Sb) and Me_2NbCl_3

	Me ₃ NbCl ₂	Me ₃ TaCl ₂	Me ₃ SbCl ₂	Me ₂ NbCl ₃
М	+1.84	+1.74	+1.51	+1.78
Cl	-0.45	-0.48	-0.55	-0.40 (ax) -0.41 (eq)
C H	$-0.08 \\ -0.08$	$-0.13 \\ -0.04$	$-0.07 \\ -0.03$	-0.05 -0.07





Fig. 7 Orbital correlation diagram for Me_3MCl_2 (M = Nb, Ta or Sb). Orbitals are correlated with respect to their principal AO character.

formed by the metals of the third transition series.¹⁹ However, the M–C e' orbital does not show a corresponding drop in energy between Nb and Ta. The metal s content of the a_1' orbitals and the relativistic stabilisation of the 6s orbital^{20,21} may well account for the changes in the orbital energies and the related changes in the ionisation energies discussed above.

Comparison with main group analogues

There appear to have been no studies published of main-group compounds of the type Me_2ECl_3 . For comparative purposes we have carried out calculations on Me_3SbCl_2 , the results of which

are reported in Table 2(c). The calculations show reasonable agreement between the calculated and experimentally determined dimensions 17 and IE 22 for this molecule.

Fig. 7 compares the orbital energies of Me₃SbCl₂ with those of its transition metal counterparts. The D_{3h} skeletal symmetry common to all these molecules permits a direct comparison between the energies for each orbital type. A striking feature of the figure is the energies of the Cl(p) orbitals; in the Sb compound these constitute the highest lying orbitals, whereas in the Nb compound they lie 1 eV lower. Both the Cl $p\pi$ and the Cl σ orbitals engage in very little mixing with other orbitals and their energies all lie very close together. However, the e' and e" orbitals of Nb and Ta are somewhat more delocalised than those of Sb and a Mulliken charge analysis (Table 4) indicates that the Cl charge is greater in Me_3SbCl_2 (-0.55) than in the transition metal compounds [-0.48 (Ta), -0.45 (Nb)]. Such a feature is also apparent in the structural parameter Δ [= r(M-Cl)_{ax} - r(M-Cl)_{eq}] observed for NbCl₅ (Δ = 0.31 Å)²³ compared with SbCl₅ ($\Delta = 0.061$ Å),²⁴ and serves to emphasise the orbitally deficient nature of the Clax-Sb-Clax unit in the main group compound.

The metal–carbon bonding levels are considerably more stable in Me₃SbCl₂ with the more electronegative Sb, in keeping with the stability of the Sb 5s and 5p orbitals. In all cases the metal–chlorine interaction is that of a 3-centre 4-electron bond. With the transition metal compounds it is the a_1' level which is M–Cl σ bonding, as the metals use a d_2 -s hybrid to form this bond and the a_2'' Cl σ level is non-bonding. For Me₃SbCl₂, by contrast, the a_2'' level is M–Cl bonding, being formed with the Sb 5p₂ orbital, and the a_1' Cl σ level is non-bonding.

Ligand site preferences

Rossi and Hoffmann⁵ concluded that the electron-rich axial sites in a d⁰ tbp system would be preferred by more electronegative ligands, and that the equatorial sites would afford the stronger π bonding opportunity. These considerations are antagonistic for a ligand such as Cl. The calculated geometries clearly indicate the preference of Cl for the axial positions. The degree of Cl to Nb π donation appears, at best, to be modest for axially located chlorine atoms. In Me2NbCl3 though, the unique equatorial chlorine enjoys a significant degree of π bonding with the suitably disposed d_{xy} and $d_{x^2-y^2}$ orbitals. The 1b₁ and 1b₂ MO's show an important contribution from Nb 5d (Table 3 and Fig. 5). Evidently the stronger M-Cl σ-interaction experienced by the axial chlorine ligands is sufficient to encourage preferential occupation of these sites despite the favourable π bonding opportunities in the equatorial positions. However, our study disagrees with the EHMO analysis of Rossi and Hoffmann⁵ in one important respect; these authors presumed the involvement of M np orbitals in the bonding of the tbp systems they considered. We find no significant population of Nb 5p from our analysis of the calculated results, in keeping with the conclusions of a recent analysis of the stereochemistries of metal alkyls and hydrides from a valence bond perspective.25

Skeletal angles in Me₂NbCl₃

A particularly striking feature of the calculated geometry of Me_2NbCl_3 is the Cl_{eq} -Nb- Cl_{ax} angle of 98.1°. Gas-phase electron diffraction studies, which will be reported elsewhere,^{26,27} are in agreement with the calculated geometry. Such a deviation from the ideal angle of 90° finds no explanation in the analysis of the PE spectrum or of the calculated MO's, there being no sign of any contribution to the 2b₁ orbital from d_{xz} on Nb (see Fig. 6). The tendency of M-Cl bonds to have large angles to adjacent bonds in early transition metal d⁰ systems, in direct contrast to their behaviour in analogous main group situations,^{1,2} is a phenomenon which has been appreciated for

certain alkyl titanium systems.^{27,28} The two most consistent explanations of this behaviour evoke either (i) Coulomb repulsion between the Cl atoms bonded to an electropositive metal centre ^{27,28} or (ii) polarization of the electron density in the (n - 1) core level of the metal.^{29,30} Given the charges calculated here (see Table 4) and the softness of the skeletal bending modes shown by the frequency calculation, we prefer the former explanation.

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